

# Lecture 36 - Bipolar Junction Transistor

(*cont.*)

December 2, 2002

## Contents:

1. Current-voltage characteristics of ideal BJT (*cont.*)
2. Charge-voltage characteristics of ideal BJT
3. Small-signal behavior of ideal BJT

## Reading material:

del Alamo, Ch. 11, §§11.2 (11.2.5), 11.3, 11.4 (11.4.1)

## Announcements:

Note special schedule for the end of the semester:

- Dec. 5: lecture
- Dec. 6: lecture
- Dec. 9: guest lecture by Prof. H. Tuller
- Dec. 11: recitation

All in regular room and regular meeting times.

## Key questions

- How do the output characteristics of the ideal BJT look like?
- How do the charge-voltage characteristics of the ideal BJT look like?
- What is the topology of the small-signal equivalent circuit model of the ideal BJT in FAR?
- What are the key dependencies of its elements?

## 1. Current-voltage characteristics of ideal BJT (*cont.*)

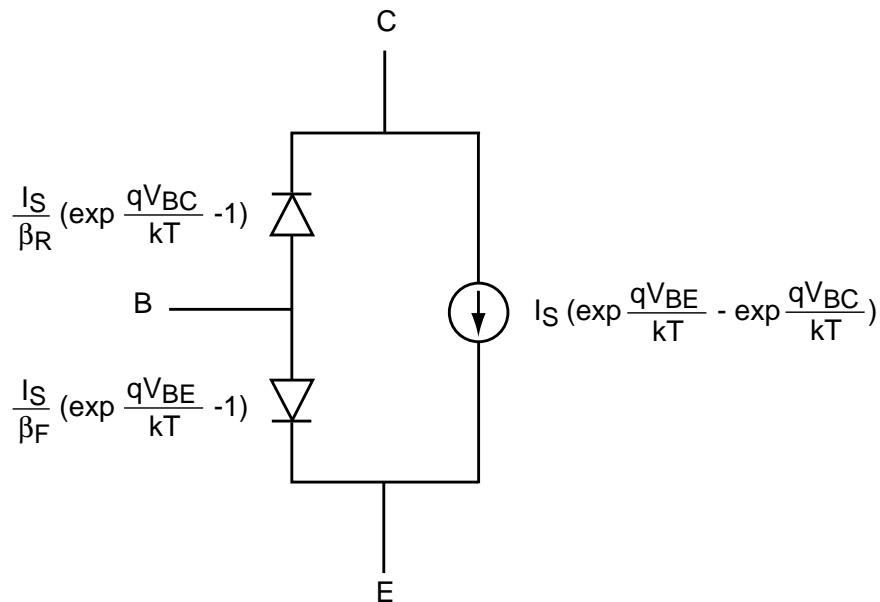
Ideal BJT current equations (superposition of forward active + reverse):

$$I_C = I_S \left( \exp \frac{qV_{BE}}{kT} - \exp \frac{qV_{BC}}{kT} \right) - \frac{I_S}{\beta_R} \left( \exp \frac{qV_{BC}}{kT} - 1 \right)$$

$$I_B = \frac{I_S}{\beta_F} \left( \exp \frac{qV_{BE}}{kT} - 1 \right) + \frac{I_S}{\beta_R} \left( \exp \frac{qV_{BC}}{kT} - 1 \right)$$

$$I_E = -\frac{I_S}{\beta_F} \left( \exp \frac{qV_{BE}}{kT} - 1 \right) - I_S \left( \exp \frac{qV_{BE}}{kT} - \exp \frac{qV_{BC}}{kT} \right)$$

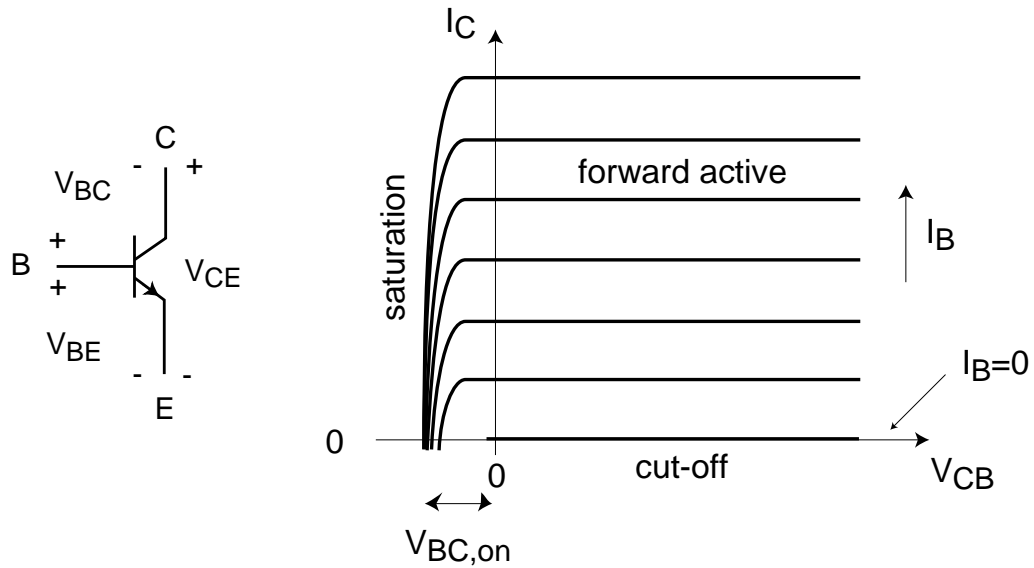
Equivalent circuit model representation:



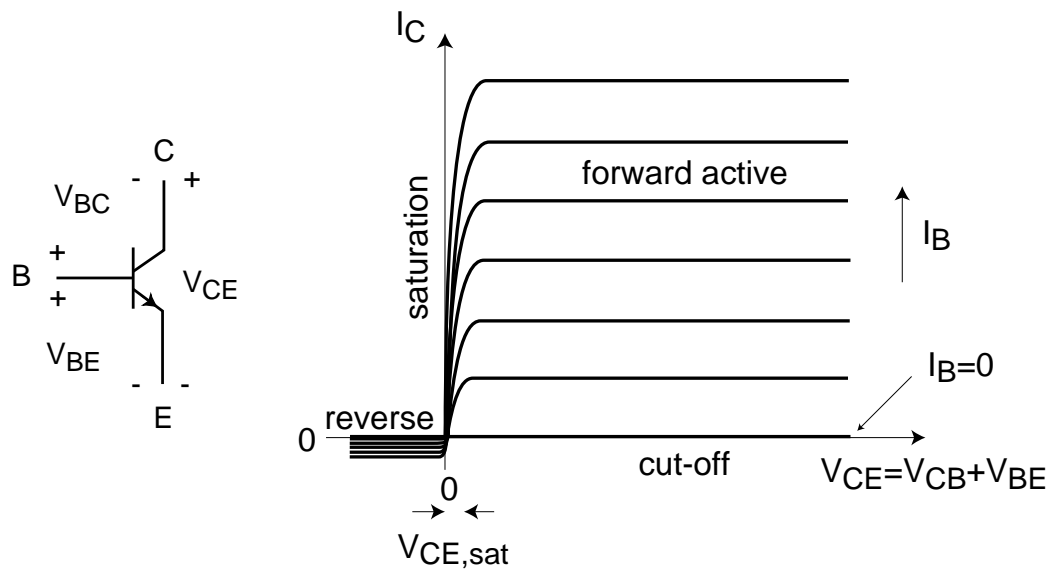
Complete model has only three parameters:  $I_S$ ,  $\beta_F$ , and  $\beta_R$ .

## □ Common-emitter output I-V characteristics

vs.  $V_{CB}$ :

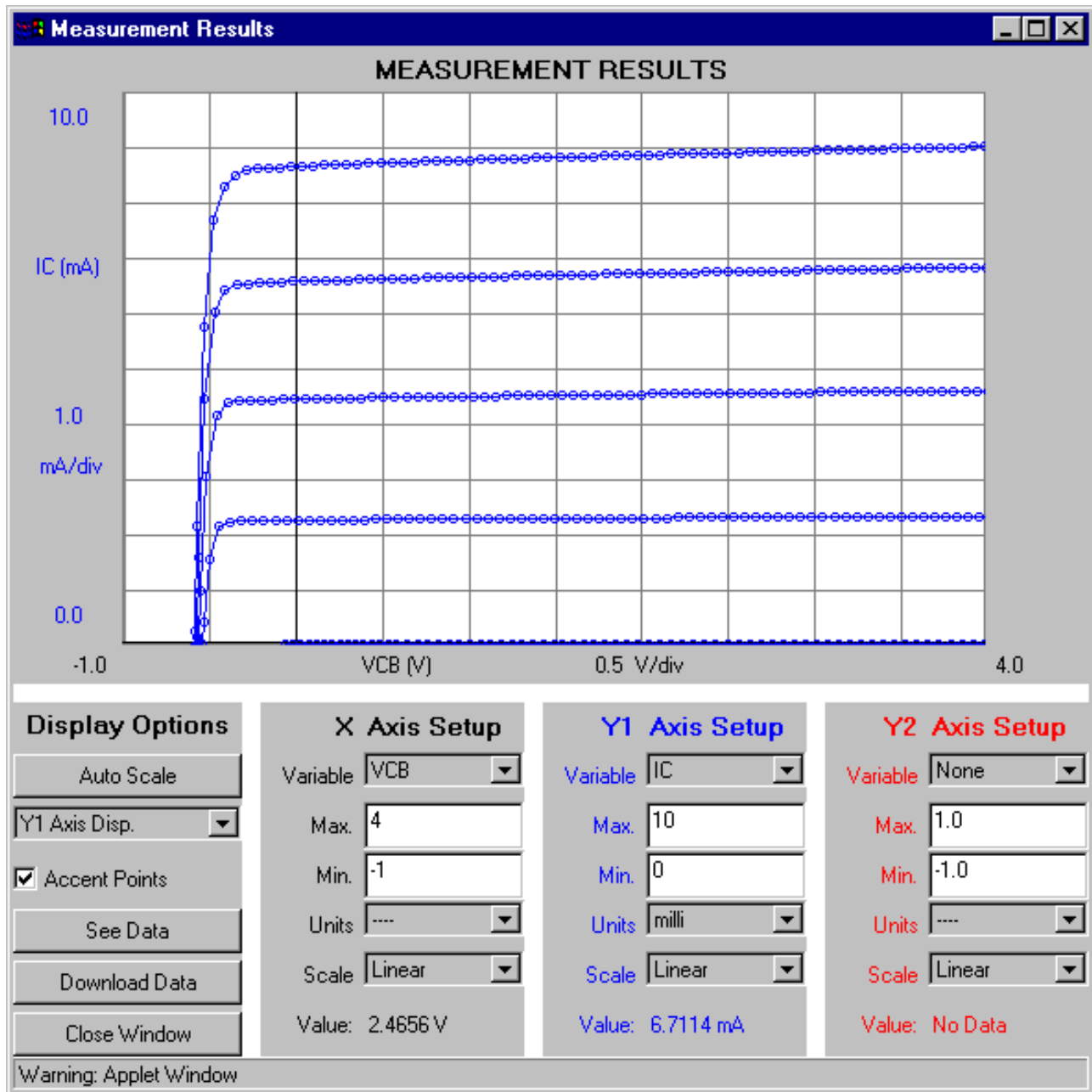


vs.  $V_{CE}$ :

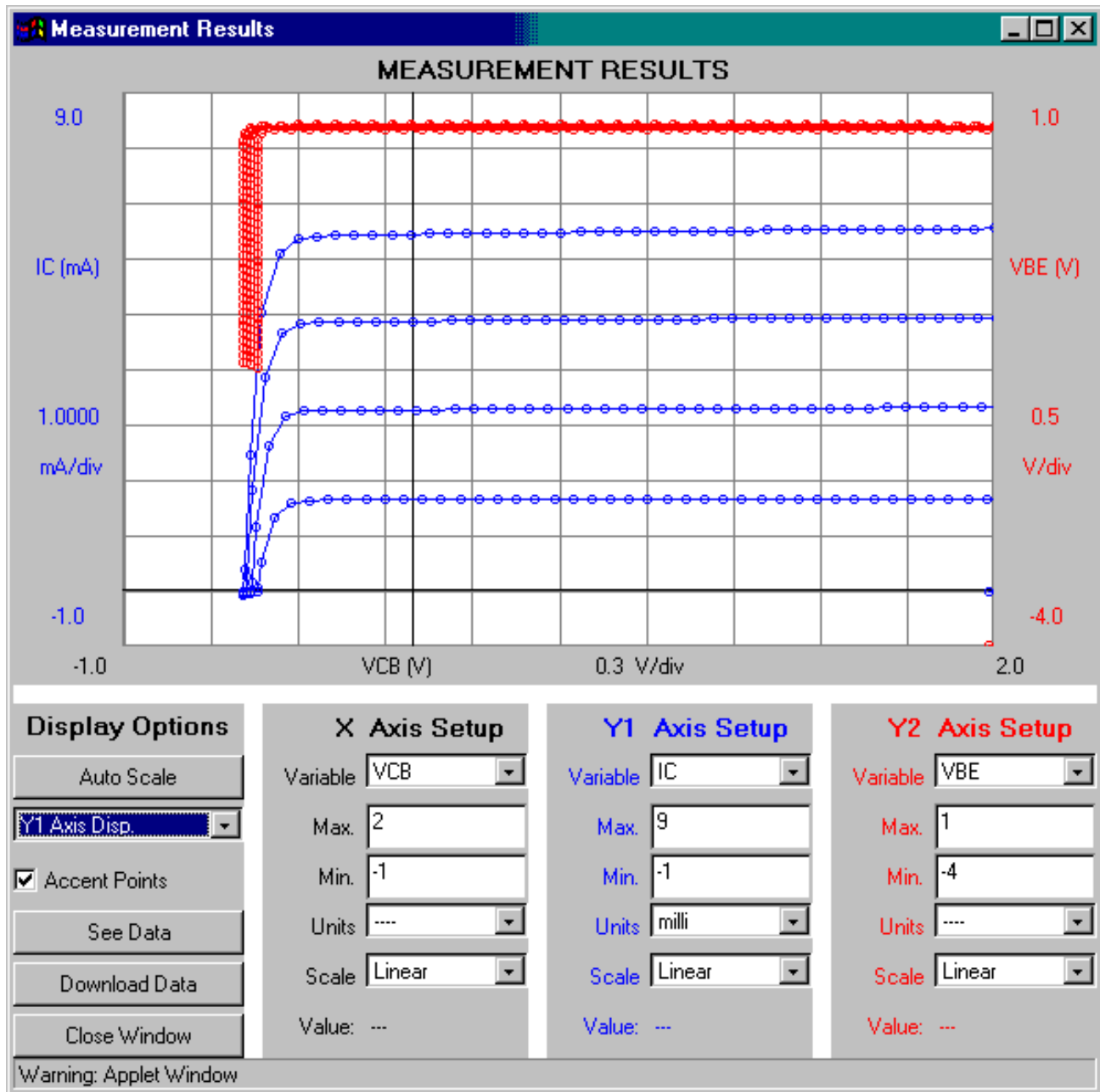


$$V_{CEsat} = -V_{BCon} + V_{BEon}$$

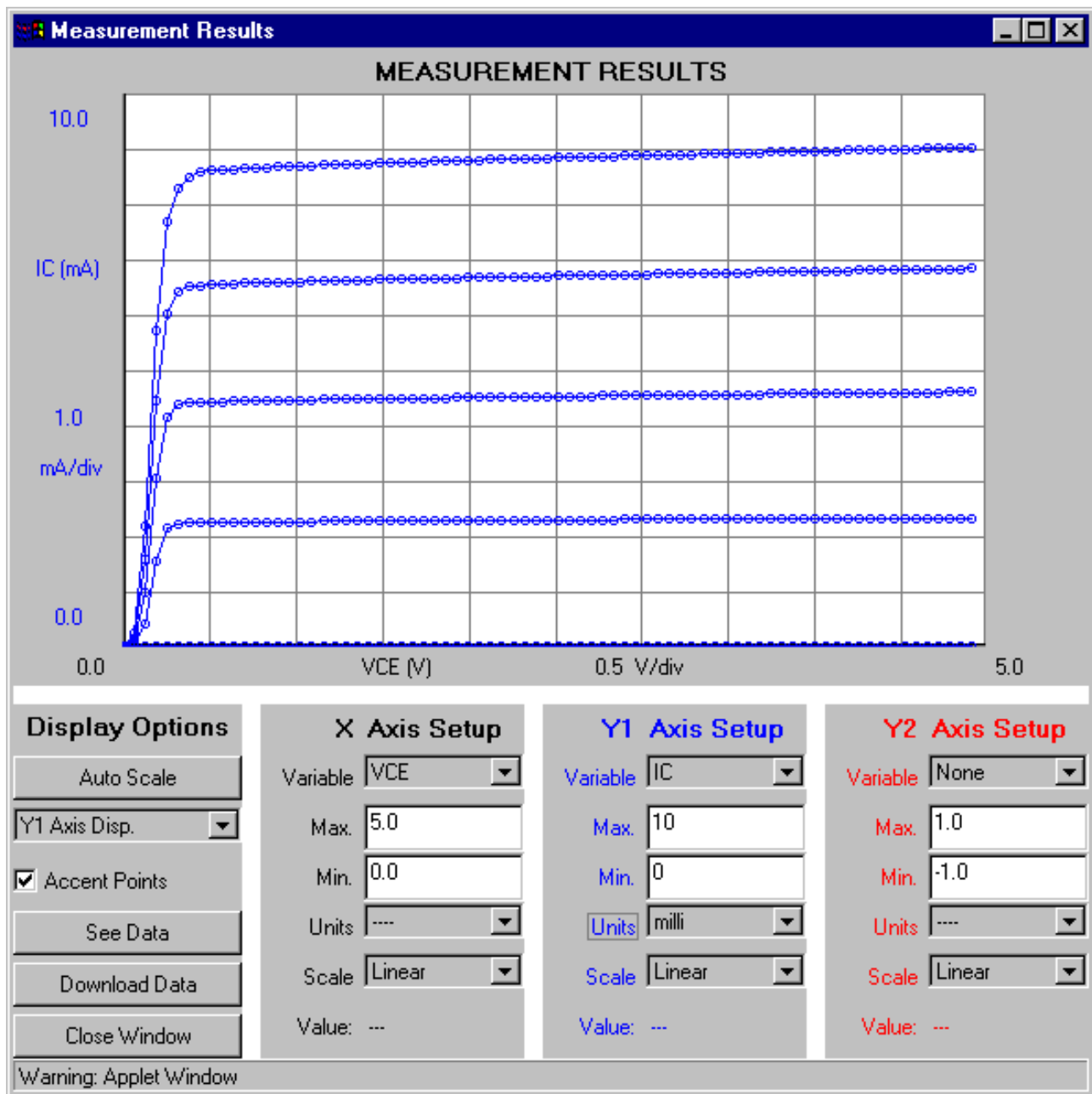
$I_C$  vs.  $V_{CB}$  with  $I_B$  as parameter:



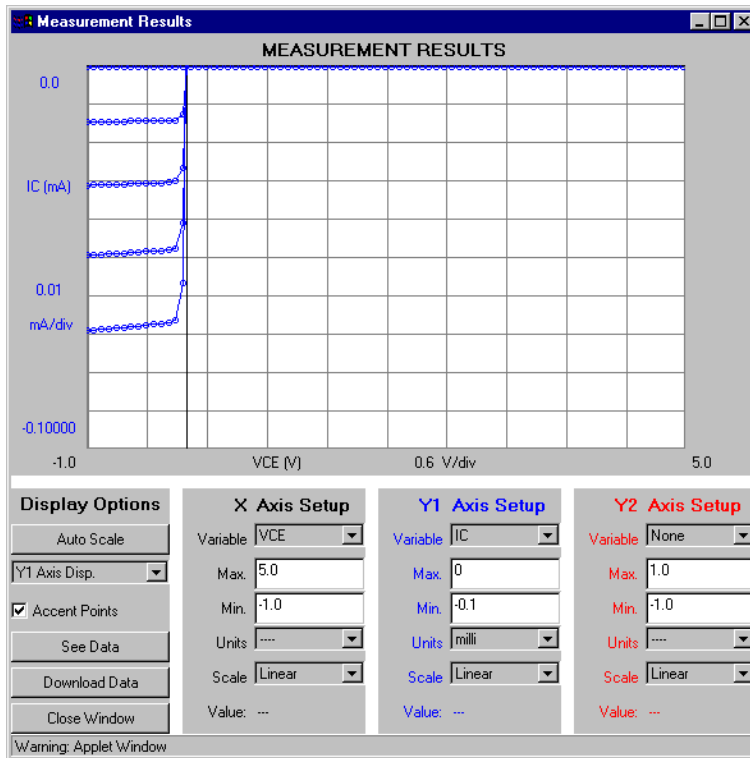
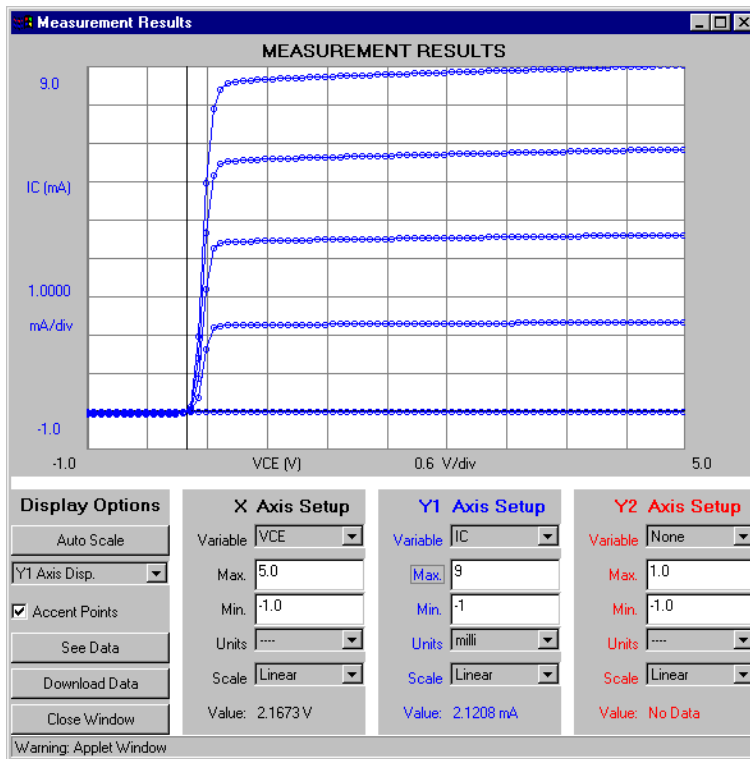
Where is the reverse regime?



Common-emitter output characteristics:



Zoom into inverse regime:



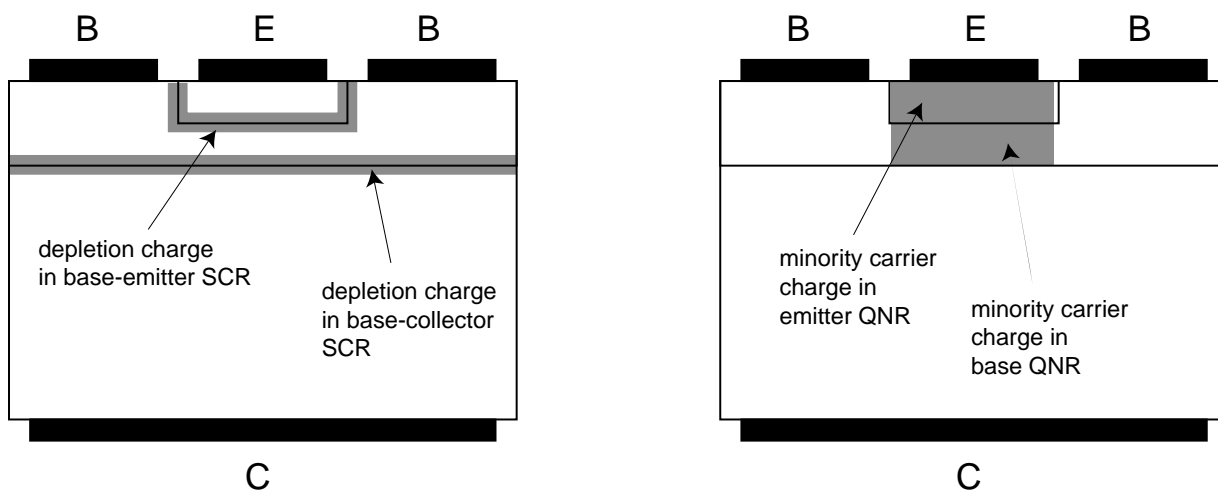


## 2. Charge-voltage characteristics of ideal BJT

In BJT, two types of stored charge:

- depletion layer charge
- minority carrier charge

In forward-active regime:



## □ Depletion layer charge

In B-E and B-C SCR's, respectively:

$$Q_{jE} = A_E \sqrt{\frac{2\epsilon q N_E N_B (\phi_{biE} - V_{BE})}{N_E + N_B}}$$

$$Q_{jC} = A_C \sqrt{\frac{2\epsilon q N_B N_C (\phi_{biC} - V_{BC})}{N_B + N_C}}$$

$\phi_{biE}$  and  $\phi_{biC}$  are respective built-in potentials.

Since  $N_E \gg N_B \gg N_C$ ,

$$Q_{jE} \simeq A_E \sqrt{2\epsilon q N_B (\phi_{biE} - V_{BE})}$$

$$Q_{jC} \simeq A_C \sqrt{2\epsilon q N_C (\phi_{biC} - V_{BC})}$$

Depletion capacitance:

$$C_{je} = \frac{\partial Q_{jE}}{\partial V_{BE}} \simeq A_E \sqrt{\frac{\epsilon q N_B}{2(\phi_{biE} - V_{BE})}} = \frac{C_{jeo}}{\sqrt{1 - \frac{V_{BE}}{\phi_{biE}}}}$$

$$C_{jc} = \frac{\partial Q_{jC}}{\partial V_{BC}} \simeq A_C \sqrt{\frac{\epsilon q N_C}{2(\phi_{biC} - V_{BC})}} = \frac{C_{jco}}{\sqrt{1 - \frac{V_{BC}}{\phi_{biC}}}}$$

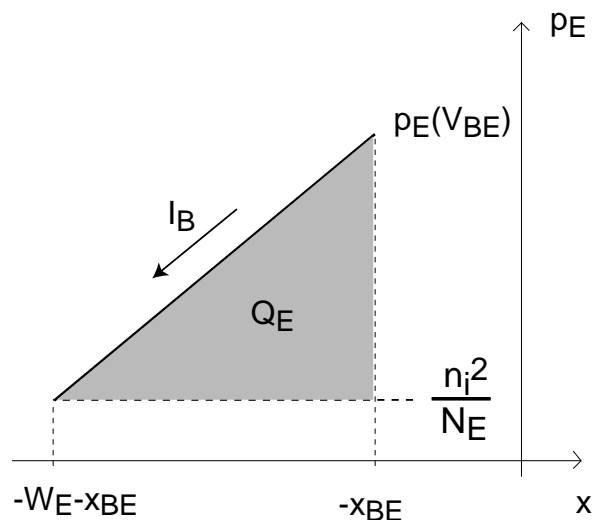
## □ Minority carrier charge

Excess minority carriers in QNR's  $\Rightarrow$  excess majority carriers to keep quasi-neutrality  $\Rightarrow$  diffusion capacitance.

Key result from pn diode: in "short" or "transparent" QNR:

$$\begin{aligned} \text{stored charge} &= \text{minority carrier transit time} \\ &\times \text{injected minority carrier current} \end{aligned}$$

- For **emitter** in FAR:

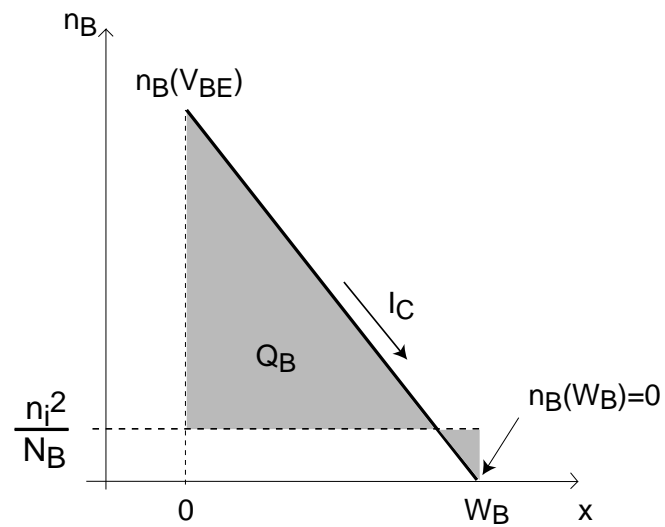


$$Q_E = \tau_{tE} I_B$$

with hole transit time:

$$\tau_{tE} = \frac{W_E^2}{2D_E}$$

- For **base** in FAR:



$$Q_B = \tau_{tB} I_C$$

with electron transit time:

$$\tau_{tB} = \frac{W_B^2}{2D_B}$$

Comments:

- Units of  $Q_E$  and  $Q_B$  are  $C$ .
- $Q_E$  and  $Q_B$  scale with  $A_E$ .

Total minority carrier charge in FAR:

$$Q_F = Q_E + Q_B = \tau_{tE}I_B + \tau_{tB}I_C = \left(\frac{\tau_{tE}}{\beta_F} + \tau_{tB}\right)I_C = \tau_F I_C$$

$$\tau_F \equiv \text{intrinsic delay [s]}$$

$\tau_F$  is overall time constant for minority carrier storage in BJT in FAR:

$$\tau_F = \frac{\tau_{tE}}{\beta_F} + \tau_{tB}$$

Note: emitter contribution to  $\tau_F$  is  $\tau_{tE}/\beta_F$  because  $I_B$  is  $\beta_F$  times smaller than  $I_C$ .

If  $V_{BE}$  changes,  $Q_E$  and  $Q_B$  change  $\Rightarrow$  capacitive effect:

$$C_F = \frac{dQ_F}{dV_{BE}} = \tau_F \frac{qI_C}{kT}$$

Location of this capacitance? Think of which terminals supply stored charge (minority and majority carriers):

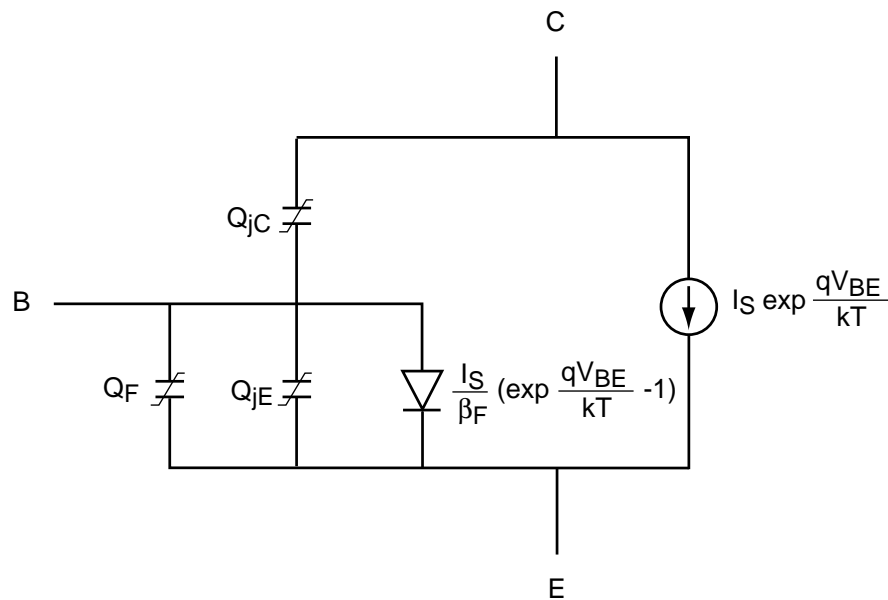
For  $Q_E$ :

- minority carriers (holes) injected from *base*
- majority carriers (electrons) come from *emitter* contact

For  $Q_B$ :

- minority carriers (electrons) injected from *emitter*
- majority carriers (holes) come from *base* contact

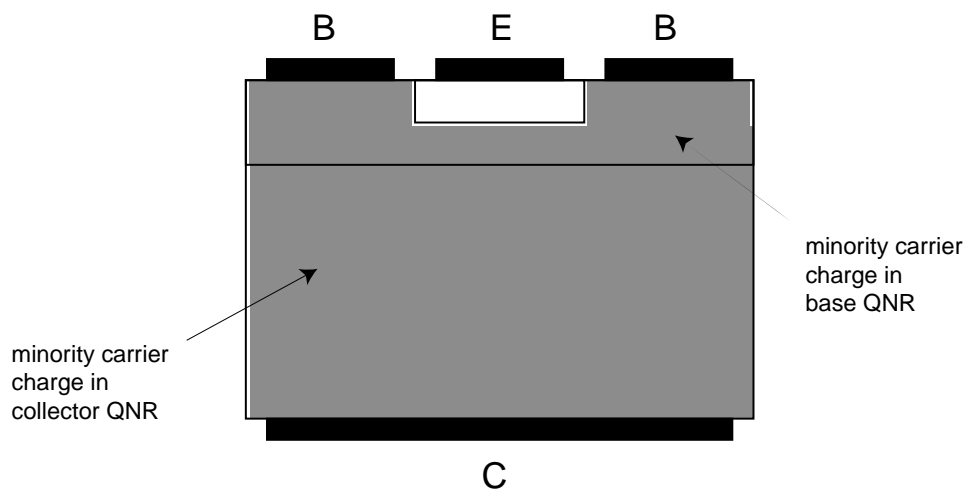
Equivalent-circuit model:



Similar picture in reverse regime: charge storage in base and collector

$$Q_R = \tau_R I_E$$

$\tau_R$  a bit complicated because it accounts for charge storage in *intrinsic* and *extrinsic* base and collector regions.

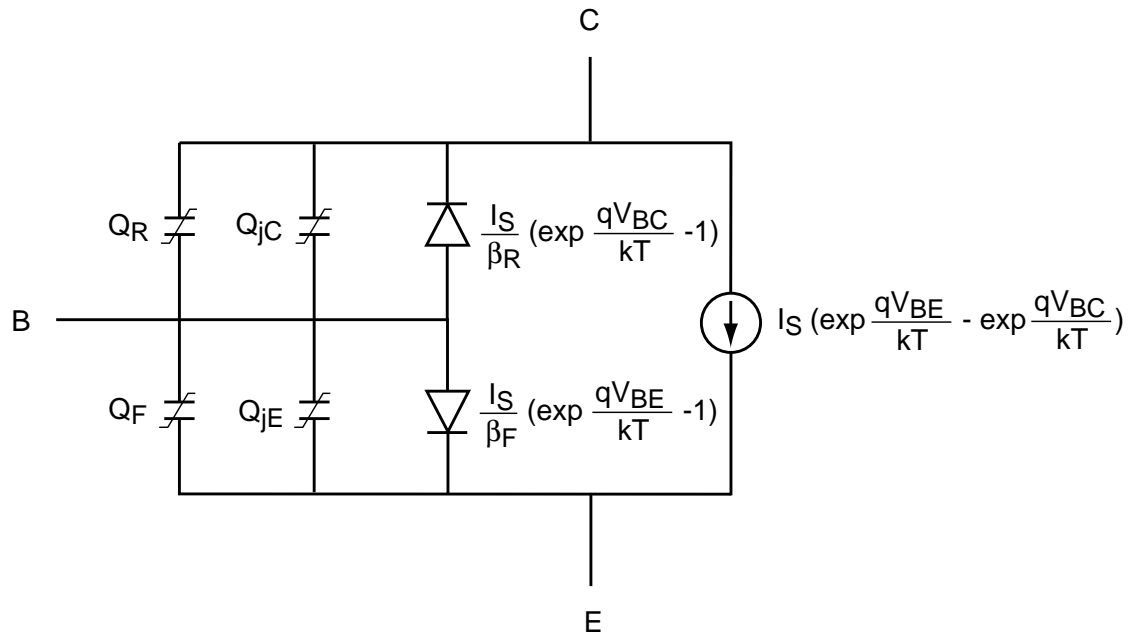


Diffusion capacitance:

$$C_R = \frac{dQ_R}{dV_{BC}} = \tau_R \frac{qI_E}{kT}$$

Located between base and collector terminals.

By superposition, complete equivalent circuit model valid in all four regimes:





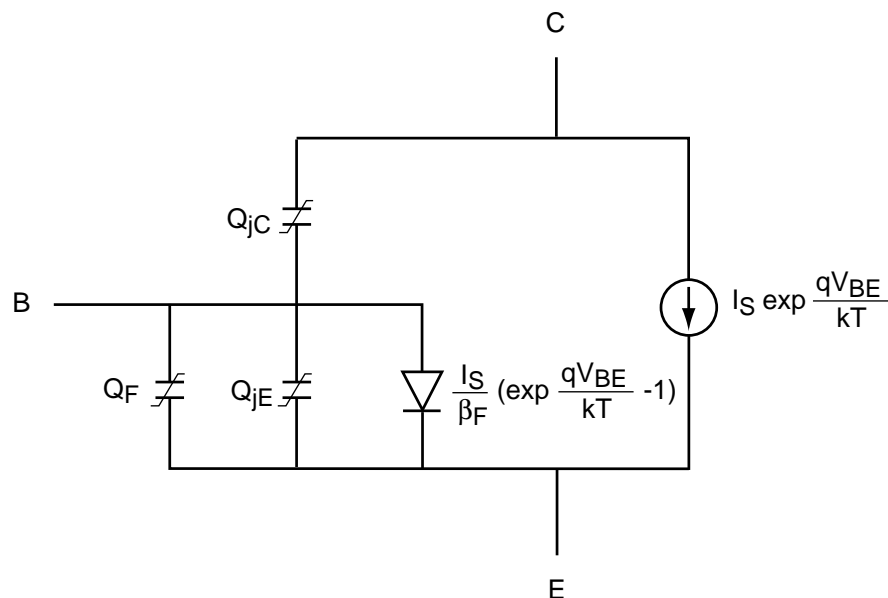
### 3. Small-signal behavior of ideal BJT

In analog (and digital) applications, interest in behavior of BJT to *small-signal* applied on top of bias

⇒ small-signal equivalent circuit model.

#### □ Small-signal equivalent circuit model in FAR

Must *linearize* hybrid- $\pi$  model in FAR:



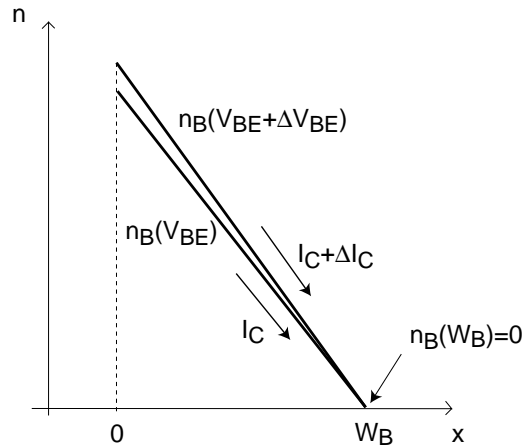
-Non-linear voltage-controlled current source linearized to *linear voltage-controlled current source*.

-Diode linearized to *resistor*.

-Charge storage elements linearized to *capacitors*.

- *Linearized voltage-controlled current source*

Apply small signal  $v_{be}$  on top of bias  $V_{BE}$ .



Collector current:

$$I_C + i_c = I_S \exp \frac{q(V_{BE} + v_{be})}{kT} \simeq I_S \exp \frac{qV_{BE}}{kT} \left(1 + \frac{qv_{be}}{kT}\right) = I_C \left(1 + \frac{qv_{be}}{kT}\right)$$

Small-signal collector current:

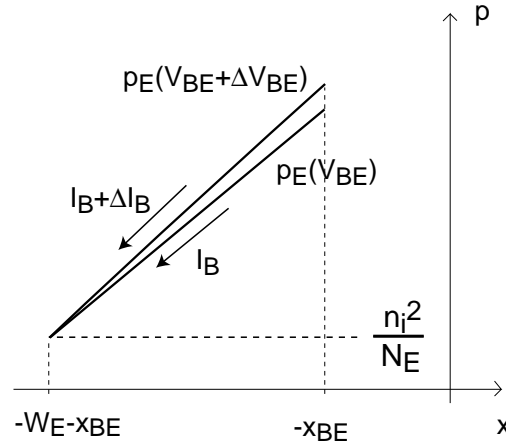
$$i_c = \frac{qI_C}{kT} v_{be}$$

Define *transconductance*:

$$g_m = \frac{qI_C}{kT}$$

$g_m$  depends only on absolute value of  $I_C$  and  $T$  (unlike MOSFET, where  $g_m$  depends on device geometry)

- *Linearized diode*



Base current:

$$I_B + i_b = I_S \exp \frac{q(V_{BE} + v_{be})}{kT} \simeq \frac{I_S}{\beta_F} \exp \frac{qV_{BE}}{kT} \left( 1 + \frac{qv_{be}}{kT} \right)$$

Small-signal base current:

$$i_b = \frac{qI_B}{kT} v_{be}$$

Define *conductance*:

$$g_\pi = \frac{qI_B}{kT} = \frac{q}{kT} \frac{I_C}{\beta_F} = \frac{g_m}{\beta_F}$$

Then, in general,

$$g_\pi \ll g_m$$

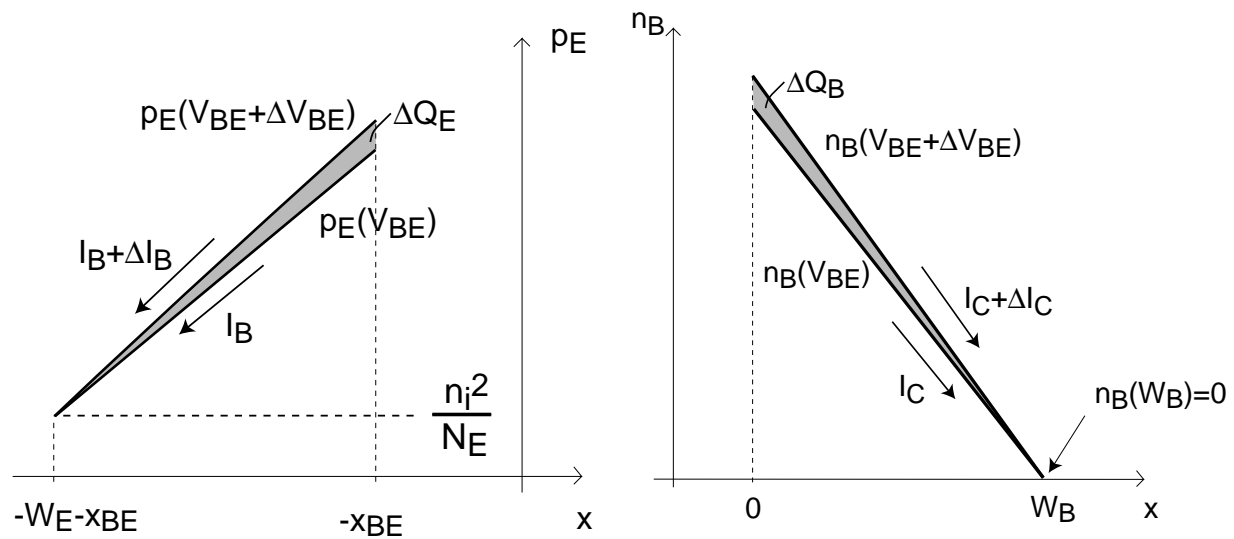
• *Capacitors*

$$Q_{jE} \rightarrow C_{je}$$

$$Q_{jC} \rightarrow C_{jc}$$

$$Q_F \rightarrow C_\pi$$

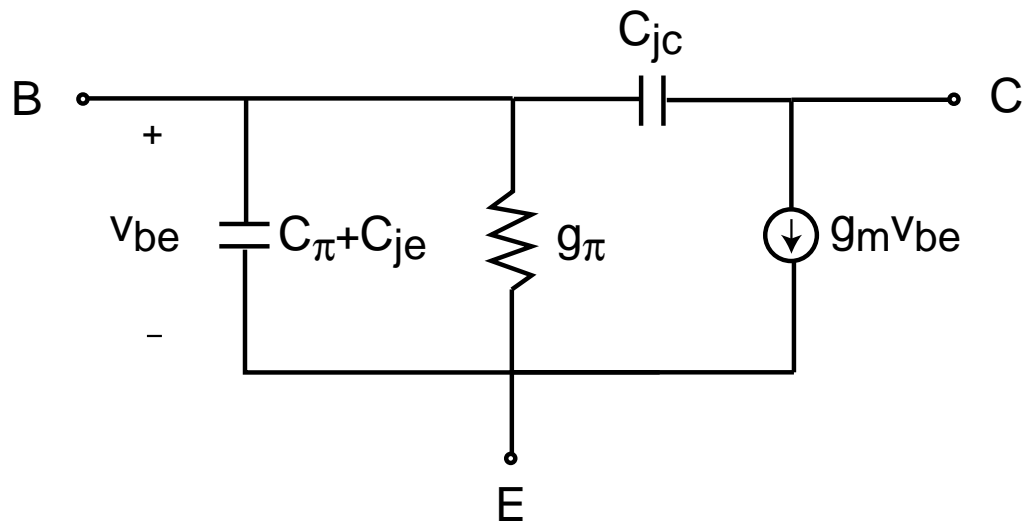
Two components in  $C_\pi$ :



Note:

$$C_\pi = \tau_F g_m$$

- Small-signal equivalent circuit model for ideal BJR in FAR:



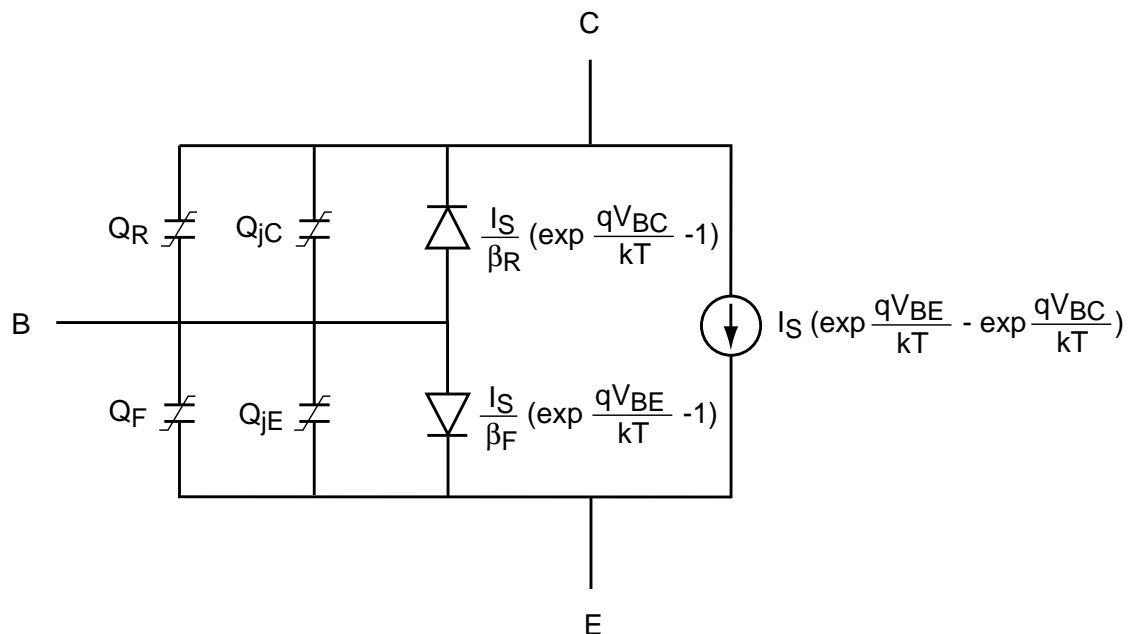
## Key conclusions

- In BJT, two types of stored charge: depletion layer charge and minority carrier charge.
- Depletion layer charge accounted through depletion capacitances.
- Minority carrier charge accounted through time constant  $\tau_F$  (*intrinsic delay*):

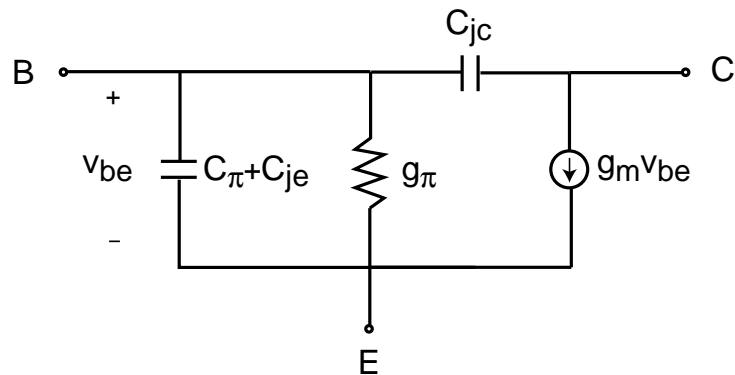
$$\tau_F = \frac{\tau_{tE}}{\beta_F} + \tau_{tB}$$

Emitter contribution to  $\tau_F$  is  $\beta_F$  times smaller than  $\tau_{tE}$  because  $I_B$  is  $\beta_F$  times smaller than  $I_C$ .

- Non-linear hybrid- $\pi$  model for ideal BJT including charge storage elements:



- Small-signal equivalent circuit model of ideal BJT in FAR:



with:

$$g_m = \frac{qI_C}{kT} \quad g_{\pi} = \frac{qI_B}{kT} = \frac{g_m}{\beta_F} \quad C_{\pi} = \tau_F g_m$$